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Institute of Mathematical Sciences
Division of Electromagnetic Research

RESEARCH REPORT No. EM-155

On an Integral Equation Arising in Inverse Scattering

C. H. YANG

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Director

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ABSTRACT

A method of Fourier transforms and double series is applied for a solution of the following integral equation arising in inverse scattering:

$$\int_{-t}^{x} R(t+y)K(x,y)dy + K(x,t) + R(x+t) = 0, t \le x$$

where R(w) is a Fourier transform of the given reflection coefficient in the differential equation $u^u(k,x)+\left[k^2-V(x)\right]u(k,x)=0$, $-\infty \le x < \infty$.

By assuming R(w) to be analytic, it is found that the unknown function can be constructed in the following forms:

i)
$$K(x,y) = \sum_{n=0}^{\infty} A_n(x) y^n$$
, ii) $K(x,y) = \sum_{p=0}^{\infty} \sum_{m+n=p} C_{m,n} x^m y^n$

in which these series converge uniformly and absolutely for |x| < x(R). x(R) depends on R(w). Therefore, if the reflection coefficient of the above differential equation is given with suitable conditions, V(x) can be found by solving the integral equation and by the relation $V(x) = 2 \frac{d}{dx} K(x,x)$. Some related topics are discussed.

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1. Introduction

Let us consider the following differential equation,

$$\frac{d^{2}}{dx^{2}} U(k,x) + [k^{2} - V(x)] U(k,x) = 0,$$

which appears in many physical problems.

If $V(\mathbf{x})$ is given, we can obtain, in principle, a solution such that

where r(k), t(k) are respectively reflection and transmission coefficients.

Conversely, one can obtain the ionization density from a knowledge of the reflection coefficient. The latter data are obtainable from the time delay of a pulsed radio wave transmitted from the earth and reflected back by the ionosphere. This ionization density is calculated as follows:

a) we first determine the kernel $K(\mathbf{x},\mathbf{y})$ from the integral equation

(1)
$$\int_{-t}^{x} R(t+y) K(x,y) dy + K(x,t) + R(x+t) = 0, t \le x$$

where R(t) = $\frac{1}{2\pi} \int_{-\infty}^{\infty} \mathbf{r}(k) e^{-i\mathbf{k}t} d\mathbf{k}$, $\mathbf{r}(k)$ being a complex reflection coefficient.

^{*}See [4] Appendix and [1] p. 5-7.

^{***}See [1], section 3. We assume that R(t) is analytic for t > 0 in this report unless specified.

- b) We next calculate $V(x) = 2 \frac{d}{dx} K(x,x)$.
- c) Finally V(x) is proportional to the ionization density.

Equation (1) can be regarded as the following Fredholm integral equation

(2)
$$-\lambda \int_{-x}^{x} R(t+y) f(y) dy + f(t) + g(t) = 0$$
 (x = fixed)

with the condition R(s) = 0 for s < 0; $\lambda = -1$, f(y) = K(x,y) and g(t) = R(x+t).

It follows that for every te[-x,x], f(t) has a solution belonging to L_2 (or C) if $R(s) \in L_2$ (or C), where L_2 is the class of square summable functions (C being the class of continuous functions).

Evaluation of a solution of (2) is very difficult, due to the condition R(s) = 0 for s < 0, by the standard Fredholm theory. It is therefore desirable to develop a new method of solving (2). From (1), after changing variables z = t+y, w = x+t, we obtain

(3)
$$\int_{0}^{w} R(z) K(x,x+z-w) dz + K(x,w-x) + R(w) = 0, w \ge 0$$

which can also be written as follows:

(4)
$$\int_{0}^{W} R(w-u) K(x,x-u) du + K(x,w-x) + R(w) = 0, w \ge 0.$$

Applying a Fourier transform to (4), for the intervals (0, ∞) and (- ∞ , 0), we obtain the following equations, respectively:

(1)
$$r_{+}(\alpha) \int_{-\infty}^{x} K(x,y) e^{-i\alpha y} dy + \int_{-x}^{\infty} K(x,y) e^{ixy} dy + e^{-i\alpha x} r_{+}(\alpha) = 0$$
for $I_{m}(\alpha) \gg 0$

See Appendix, p. 26. $I_m(\alpha)$ means the imaginary part of α ; A >> 0 means A is positive and very large. Similarly A << 0 means A is negative and very large.

$$(\text{ii}) \quad -r_{\underline{}}(\alpha) \int_{0}^{\infty} K(x,y) \, \mathrm{e}^{-\mathrm{i}\alpha y} \mathrm{d}y \, + \int_{-\infty}^{-X} K(x,y) \, \mathrm{e}^{\mathrm{i}\alpha y} \mathrm{d}y \, + \, \mathrm{e}^{-\mathrm{i}\alpha x} r_{\underline{}}(\alpha) \, = \, 0$$

$$\qquad \qquad \qquad \text{for } I_{m}(\alpha) \, < < \, 0$$

where
$$r_{+}(\alpha) = \int_{0}^{\infty} R(t) e^{i\alpha t} dt = r(\alpha)$$
, $r_{-}(\alpha) = \int_{\infty}^{0} R(t) e^{i\alpha t} dt$

Here, we have extended the definition of R(t) for t < 0.

Substituting - α for α into i) and ii), we obtain, respectively, the following:

$$\begin{array}{lll} \mbox{(iii)} & & r_{+}(-\alpha) \, \int_{-\infty}^{X} \! K(x,\!y) \, \mathrm{e}^{\, \mathrm{i} \, \alpha \! y} \mathrm{d} y \, + \, \int_{-X}^{\infty} \! K(x,\!y) \, \mathrm{e}^{\, -\mathrm{i} \, \alpha \! y} \mathrm{d} y \, + \, \mathrm{e}^{\, \mathrm{i} \, \alpha \! x} \, r_{+}(-\alpha) \, = \, 0 \\ & & \mbox{for } \, \mathrm{I}_{\mathrm{m}}(\alpha) \, < < \, 0 \\ \\ \mbox{(iv)} & & -r_{-}(-\alpha) \, \int_{X}^{\infty} \! K(x,\!y) \, \mathrm{e}^{\, \mathrm{i} \, \alpha \! y} \mathrm{d} y \, + \, \int_{-\infty}^{-X} \! K(x,\!y) \, \mathrm{e}^{\, -\mathrm{i} \, \alpha \! y} \mathrm{d} y \, + \, \mathrm{e}^{\, \mathrm{i} \, \alpha \! x} \, r_{-}(-\alpha) \, = \, 0 \\ & & \mbox{for } \, \mathrm{I}_{\mathrm{m}}(\alpha) \, > > \, 0. \end{array}$$

The sets of equations (i,iv) and (ii,iii) cannot be solved uniquely unless we have more conditions on K(x,y) since there are four unknowns but only two equations, respectively. Nevertheless, when x=0, they are solvable for K(0,y) and we obtain the following equations from (i-iv):

$$\begin{split} &(\text{i'}) \qquad r_{+}(\alpha) \int_{-\infty}^{O} K(\text{O},y) \, \mathrm{e}^{-\mathrm{i}\alpha y} \mathrm{d}y \, + \int_{O}^{\infty} \!\! K(\text{O},y) \, \mathrm{e}^{\mathrm{i}\alpha y} \mathrm{d}y \, + \, r_{+}(\alpha) \, = \, 0 \, , \, \, \mathrm{I}_{\mathrm{m}}(\alpha) \, \gg \, 0 \\ &(\text{ii'}) \quad -\mathrm{r}_{-}(\alpha) \int_{O}^{\infty} \!\! K(\text{O},y) \, \, \mathrm{e}^{-\mathrm{i}\alpha y} \mathrm{d}y \, + \, \int_{-\infty}^{O} \!\! K(\text{O},y) \, \mathrm{e}^{\mathrm{i}\alpha y} \mathrm{d}y \, + \, r_{-}(\alpha) \, = \, 0 \, , \, \, \mathrm{I}_{\mathrm{m}}(\alpha) \, \ll \, 0 \\ &(\text{iii'}) \quad r_{+}(-\alpha) \int_{-\infty}^{O} K(\text{O},y) \, \mathrm{e}^{\mathrm{i}\alpha y} \mathrm{d}y \, + \, \int_{O}^{\infty} \!\! K(\text{O},y) \, \mathrm{e}^{-\mathrm{i}\alpha y} \mathrm{d}y \, + \, r_{+}(-\alpha) \, = \, 0 \, , \, \, \, \mathrm{I}_{\mathrm{m}}(\alpha) \, \ll \, 0 \end{split}$$

×

We take the analytical continuation of R(t) as its extended definition for t \leq 0, in the above integral.

From (i') and (iv') we obtain

(5)
$$\hat{K}_{+}(0,\alpha) = \int_{0}^{\infty} K(0,y) e^{i\alpha y} dy = \begin{vmatrix} r_{+}(\alpha), r_{+}(\alpha) \\ 1, r_{-}(-\alpha) \end{vmatrix} / \begin{vmatrix} 1, r_{+}(\alpha) \\ -r_{-}(-\alpha), 1 \end{vmatrix}$$
$$= r_{+}(\alpha) \left[r_{-}(-\alpha) - 1\right] / \left[1 + r_{+}(\alpha)r_{-}(-\alpha)\right].$$

Similarly, from (ii') and (iii'), we obtain

(6)
$$\hat{K}_{-}(0,\alpha) = \int_{-\infty}^{0} K(0,y) e^{i\alpha y} dy = -r_{-}(\alpha) [1 + r_{+}(-\alpha)] / [1 + r_{+}(-\alpha) r_{-}(\alpha)].$$

It follows that

(7)
$$K(0,y) = \frac{1}{2\pi} \int_{1c-\infty}^{1c+\infty} \hat{K}_{+}(0,\alpha) e^{-i\alpha y} dy + \frac{1}{2\pi} \int_{1d-\infty}^{1d+\infty} \hat{K}_{-}(0,\alpha) e^{-i\alpha y} dy$$

for $c>\!\!>0$ and $d<\!\!<0$.

According to Fredholm's theory, the solution of (2) can be written as follows:

$$f(t) = -g(t) - \int_{-x}^{x} K_{R}(t,y;\lambda)g(y)dy$$

where

$$\begin{split} & K_{R}(\textbf{t},\textbf{y}) = \textbf{D}(\textbf{t},\textbf{y};\boldsymbol{\lambda}) \middle/ \textbf{D}(\boldsymbol{\lambda}) & \text{if } \textbf{D}(\boldsymbol{\lambda}) \neq \textbf{0} \text{ ,} \\ & \textbf{D}(\boldsymbol{\lambda}) = \textbf{1} + \sum_{n=1}^{\infty} \frac{(-\lambda)^{n}}{n!} \int_{-x}^{x} \dots \int_{-x}^{x} \textbf{R} \begin{pmatrix} \textbf{s}_{1},\dots,\textbf{s}_{n} \\ \textbf{s}_{1},\dots,\textbf{s}_{n} \end{pmatrix} d\textbf{s}_{1} \dots d\textbf{s}_{n} \text{ ,} \\ & \textbf{D}(\textbf{t},\textbf{y};\boldsymbol{\lambda}) = \textbf{1} + \sum_{n=1}^{\infty} \frac{(-\lambda)^{n}}{n!} \int_{-x}^{x} \dots \int_{-x}^{x} \textbf{R} \begin{pmatrix} \textbf{t},\textbf{s}_{1},\dots,\textbf{s}_{n} \\ \textbf{y},\textbf{s}_{1},\dots,\textbf{s}_{n} \end{pmatrix} d\textbf{s}_{1} \dots d\textbf{s}_{n} \text{ ,} \end{split}$$

and
$$R\begin{pmatrix} x_1, \dots, x_n \\ y_1, \dots, y_n \end{pmatrix} = \begin{pmatrix} R(x_1 + y_1), \dots, R(x_1 + y_n) \\ - - - - - - - - \\ R(x_n + y_1), \dots, R(x_n + y_n) \end{pmatrix}$$

It follows that the solution of (1)

$$K(x,t) = -R(x+t) - \frac{1}{D(-1)} \int_{-1}^{x} D(t,y;-1) R(x+y) dy$$

is a function analytic with respect to t and meromorphic with respect to x if R(w) is analytic with respect to w.

If $|\lambda_O(x)| > 1$ then D(-1) \neq 0, where $\lambda_O(x)$ is the first eigenvalue of (2), and this is always true if x is sufficiently small since $|\lambda_O(x)|^2$ is greater than $\left\{\int_0^x \int_{-y}^x |R(t+y)|^2 dt \ dy\right\}^{-1}.$

It follows that K(x,t) can always be expanded in a power series in x and t for small x. And K(x,y) can also be expanded in a power series in x with coefficients as functions of y. In the following section, we shall discuss the two cases in which we could find the solution of (1) by a method of Fourier transforms.

2. The case where $K(x,y) = \sum_{n=0}^{\infty} \frac{A_n(y)x^n}{n}$

As in Section 1, we can assume that K(x,y) can be expanded in power series of x: $K(x,y) = \sum_{n=0}^{\infty} A_n(y) x^n$ which is uniformly and absolutely convergent for small x, e.g. we can take any value of x

which is less than $x_0 = \sup \left\{ a \left| \int_0^a \int_{-y}^a \left| R(t+y) \right|^2 dt dy < 1 \right\} \right\}$.

Theorem 1:

Let R(t) be a given analytic function, then the integral equation (1) has a solution K(x,y) = $\sum_{n=0}^{\infty} A_n(y) x^n$ for $|x| < x_0$ where $A_n(y)$ can be obtained from the following recursion formula:

$$A_{n}(y) = \frac{1}{2\pi} \int_{1c-\infty}^{1c+\infty} \frac{r_{+}^{\circ}(\alpha)r_{-}^{n}(-\alpha) - r_{+}^{n}(\alpha)}{1 + r_{-}^{\circ}(\alpha)r_{-}^{\circ}(-\alpha)} e^{-iy\alpha} d\alpha \quad (c \gg 0)$$

(8)
$$\frac{A_{n}(y) = \frac{1}{2\pi} \int_{1c-\infty}^{c} \frac{1 + r_{+}^{\circ}(\alpha) r_{-}^{\circ}(-\alpha)}{1 + r_{+}^{\circ}(\alpha) r_{-}^{\circ}(-\alpha)} e^{-i\alpha} (c > 0) }{1 + r_{+}^{\circ}(\alpha) r_{-}^{\circ}(\alpha) r_{-}^{\circ}(\alpha)} e^{-iy\alpha} d\alpha \quad (d << 0)$$

where

$$(9) \hspace{1cm} r_{+}^{\hspace{1cm} n} (\alpha) \hspace{1cm} = \hspace{1cm} \int_{0}^{\hspace{1cm} \infty} R^{\hspace{-1pt} n} (0,t) e^{\hspace{-1pt} 1} dt \hspace{1cm} , \hspace{1cm} I_{\hspace{-1pt} m} (\alpha) >\!\!> 0$$

(10)
$$r_{\underline{}}^{n}(\alpha) = \int_{\underline{}}^{0} R^{n}(0,t)e^{i\alpha t}dt$$
, $I_{\underline{m}}(\alpha) \ll 0$

(11)
$$R^{n}(x,t) = \frac{1}{x} \left\{ \int_{-t}^{x} R(t+y) \mathbb{H}_{n-1}(0,y) dy + \mathbb{H}_{n-1}(0,t) + R^{n-1}(x,t) \right\}$$

(12)
$$H_n(x,t) = \frac{1}{x} \left\{ H_{n-1}(x,y) - H_{n-1}(0,y) \right\}$$
 $(n \ge 1)$

(13)
$$R^{O}(x,y) = R(x+y)$$

(14)
$$H_{O}(x,y) = K(x,y)$$
.

It is known that when n=0, $A_{O}(y) = K(O,y)$ and formula (8) becomes

identical with (7); when n=1, $A_1(y) = H_1(0,y) = \frac{\partial}{\partial x} K(x,y)$, and x=0

so on, and in general
$$A_n(y) = H_n(0,y) = \left(\frac{\partial}{\partial x}\right)^n K(x,y)$$
 .

Proof of Theorem 1:

When n=0, we have already proven the theorem in Section 1. By Mathematical Induction, suppose it has been proven for the cases k < n, then for k=n, we would have from (1)

$$(15) \qquad \int_{-t}^{x} R(t+y)H_{n}(x,y)dy + H_{n}(x,t) + R^{n}(x,t) = 0 .$$

This is true since, by our assumption, we already know that

(16)
$$\int_{-t}^{x} R(t+y) H_{n-1}(x,y) dy + H_{n-1}(x,t) + R^{n-1}(x,t) = 0$$

holds. It follows from (16) that

$$\int_{-t}^{x} R(t+y) \left\{ H_{n-1}(x,y) - H_{n-1}(0,y) \right\} dy + H_{n-1}(x,t) - H_{n-1}(0,t)$$

$$+ R^{n-1}(x,t) + \int_{-t}^{x} R(t+y) H_{n-1}(0,y) dy + H_{n-1}(0,t) = 0.$$

Dividing (17) by x, and using the formulas (11) and (12), we obtain the integral equation (15). Applying a method similar to that used in Section 1, we obtain

(I)
$$\mathbf{r}_{+}^{\circ}(\alpha) \int_{-\infty}^{x} \mathbf{H}_{n}(x,y) e^{-i\alpha y} dy + \int_{-x}^{\infty} \mathbf{H}_{n}(x,y) e^{i\alpha y} dy + \mathbf{r}_{+}^{n}(x,\alpha) = 0$$

$$\mathbf{Im}(\alpha) \gg 0$$

$$(II) \qquad -r_{}^{\circ}(\alpha) \int_{x}^{\infty} H_{n}(x,y) e^{-i\alpha y} dy + \int_{-\infty}^{-x} H_{n}(x,y) e^{i\alpha y} dy + r_{}^{n}(x,\alpha) = 0$$

$$Im(\alpha) << 0$$

where

$$\begin{split} \mathbf{r}_{+}^{n}\left(\mathbf{x,}\alpha\right) &= \int_{-\mathbf{x}}^{\infty} \mathbf{R}^{n}(\mathbf{x,}\mathbf{y}) \, \mathrm{e}^{\mathrm{i}\alpha\mathbf{y}} \mathrm{d}\mathbf{y} \\ \\ \mathbf{r}_{-}^{n}\left(\mathbf{x,}\alpha\right) &= \int_{-\infty}^{-\mathbf{x}} \mathbf{R}^{n}(\mathbf{x,}\mathbf{y}) \, \mathrm{e}^{\mathrm{i}\alpha\mathbf{y}} \mathrm{d}\mathbf{y} \end{split} .$$

Putting x=0, and using the same procedure as in Section 1, we obtain formula (8), where $r_+^n(\alpha) \equiv r_+^n(0,\alpha)$, $r_-^n(\alpha) \equiv r_-^n(0,\alpha)$.

3. The case where
$$K(x,y) = \sum_{n=0}^{\infty} \sum_{m+n=p} c_{m,n} x^m y^n$$

Theorem 2:

Let
$$R(w) = \sum_{n=0}^{\infty} \gamma_n w^n$$
 be given, then the integral equation (1)

has a solution $K(\textbf{x},\textbf{y}) = \sum_{p=0}^{\infty} \sum_{\textbf{m}+\textbf{n}=p} \textbf{C}_{\textbf{m},\textbf{n}} \textbf{x}^{\textbf{m}} \textbf{y}^{\textbf{n}} \text{ for } |\textbf{y}| \leq |\textbf{x}| \leq \textbf{x}_{\textbf{o}}, \text{ where } \textbf{x} = \textbf{x}_{\textbf{o}}$

$$c_{m,n}$$
 can be determined from γ_k by the following (m,n=0,1,2,...)

recursion formula:

$$\text{(Here, } x_{0} = \sup \left\{ x \mid \int_{0}^{x} \int_{-y}^{x} \left| R(t+y) \right|^{2} dt \ dy < 1 \right\})$$

$$\text{(18)} \quad \delta_{0,q} \gamma_{\ell} + \sum_{k=0}^{\ell-1} \sum_{m=0}^{q} \left(-1\right)^{k} \begin{pmatrix} k+q-m \\ q-m \end{pmatrix} \frac{C_{m,k+q-m} \gamma_{\ell-k-1}}{\ell \binom{\ell-1}{k}} + \sum_{m=0}^{q} \left(-1\right)^{q-m} C_{m,\ell+q-m} \begin{pmatrix} \ell+q-m \\ q-m \end{pmatrix} = 0$$

for q,l=0,1,2,3,... (Here, we use the convention $\sum_{k=0}^{-1} f_k = 0$, δ_{ij} being Kronecker's delta).

The procedure to obtain $C_{m,n}$ from (18) is as follows: Putting q=0, t=0,1,2,3..., we can obtain $C_{0,m}$ (m=0,1,2,...), step by step; then we obtain $C_{1,0}$ from $C_{0,0}$ and $C_{0,1}$, and so on. We obtain $C_{1,m}$ from $C_{0,p}$ (0 \leq p \leq m+1) and $C_{1,q}$ (0 \leq q \leq m-1). Similarly, we obtain $C_{m,n}$ from $C_{m,n}$ from $C_{m,n}$ from $C_{m,n}$ (0 \leq p \leq m-1, 0 \leq q \leq m+1) and $C_{m,n}$ (0 \leq t \leq n-1).

Proof of Theorem 2:

Since
$$K(\mathbf{x},\mathbf{y}) = \sum_{p=0}^{\infty} \sum_{m+n=p}^{\infty} \mathbf{c}_{m,n} \mathbf{x}^{m,p}$$
, we obtain
$$\begin{split} K(\mathbf{x},\mathbf{x}-\mathbf{u}) &= \sum_{p=0}^{\infty} \sum_{m+n=p}^{\infty} \mathbf{c}_{m,n} \mathbf{x}^{m} (\mathbf{x}-\mathbf{u})^{n} \\ &= \sum_{p=0}^{\infty} \sum_{m=0}^{p} \mathbf{c}_{m,p-m} \sum_{q=m}^{p} (-1)^{p-q} \begin{pmatrix} \mathbf{p}^{-m} \\ \mathbf{q}^{-m} \end{pmatrix} \mathbf{u}^{p-q} \mathbf{x}^{q} \\ &= \sum_{p=0}^{\infty} \sum_{q=0}^{p} \begin{pmatrix} \sum_{m=0}^{q} \mathbf{c}_{m,p-m} (-1)^{p-q} \begin{pmatrix} \mathbf{p}^{-m} \\ \mathbf{q}^{-m} \end{pmatrix} \mathbf{u}^{p-q} \mathbf{x}^{q} \\ &= \sum_{q=0}^{\infty} \begin{pmatrix} \sum_{p=q}^{\infty} \sum_{m=0}^{q} \mathbf{c}_{m,p-m} (-1)^{p-q} \begin{pmatrix} \mathbf{p}^{-m} \\ \mathbf{q}^{-m} \end{pmatrix} \mathbf{u}^{p-q} \end{pmatrix} \mathbf{x}^{q} \end{split}$$

Similarly

(19)
$$K(x,w-x) = \sum_{p=0}^{\infty} \sum_{m+n=p} C_{m,n} x^{m} (w-x)^{n}$$

$$= \sum_{q=0}^{\infty} \left(\sum_{p=q}^{\infty} \sum_{m=0}^{q} C_{m,p-m} (-1)^{q-m} {p-m \choose q-m} w^{p-q} \right) x^{q} ,$$

and

$$(20) \int_{0}^{W} R(w-u)K(x,x-u) du = \int_{0}^{W} R(w-u) \left\{ \sum_{q=0}^{\infty} \left(\sum_{p=q}^{\infty} \sum_{m=0}^{q} C_{m,p-m}(-1)^{p-q} \binom{p-m}{q-m} u^{p-q} \right) x^{q} \right\} du$$

$$= \sum_{q=0}^{\infty} \left\{ \int_{0}^{W} R(w-u) \left[\sum_{p=q}^{\infty} \sum_{m=0}^{q} C_{m,p-m}(-1)^{p-q} \binom{p-m}{q-m} u^{p-q} \right] du \right\} x^{q} .$$

Substituting (19) and (20) into (4) and arranging in powers of x, we obtain from the coefficient of $q^{\rm th}$ power the following:

(21)

$$+ \sum_{p=q}^{\infty} \sum_{m=0}^{q} C_{m,p-m} (-1)^{q-m} \binom{p-m}{q-m} w^{p-q} + \delta_{0,q} R(w) = 0.$$

Also, we have

$$\begin{split} & \int_{0}^{W} R(w-u) \left\{ \sum_{p=q}^{\infty} \sum_{m=0}^{q} C_{m,p-m} (-1)^{p-q} \begin{pmatrix} p-m \\ q-m \end{pmatrix} u^{p-q} \right\} du \\ & = \sum_{p=q}^{\infty} \sum_{m=0}^{q} C_{m,p-m} (-1)^{p-q} \begin{pmatrix} p-m \\ q-m \end{pmatrix} \int_{0}^{W} u^{p-q} \sum_{n=0}^{\infty} \gamma_{n} (w-u)^{n} du \end{split}$$

$$= \sum_{p=q}^{\infty} \sum_{m=0}^{q} C_{m,p-m} (-1)^{p-q} \begin{pmatrix} p-m \\ q-m \end{pmatrix} \sum_{\ell=p-q}^{\infty} \gamma_{\ell+q-p} \sum_{s=0}^{\ell+q-p} \begin{pmatrix} \ell+q-p \\ s \end{pmatrix} \frac{(-1)^{\ell+q+p+s} w^{\ell+1}}{\ell-s-1}$$

(put

$$\begin{array}{c} k = p - q) \\ = \sum_{\ell = 0}^{\infty} \left\{ \sum_{k = 0}^{\ell} \sum_{m = 0}^{q} C_{m, k + q - m} (-1)^{k} \begin{pmatrix} k + q - m \\ q - m \end{pmatrix} \cdot \frac{\gamma_{\ell - k}}{(\ell + 1) \cdot \binom{\ell}{\nu}} \right\} w^{\ell + 1} \end{array} .$$

(Here we use (*)
$$\dots \sum_{s=0}^{\ell-k} \binom{\ell-k}{s} \frac{(-1)^{\ell+k+s}}{\ell-s+1} = \frac{1}{(\ell+1) \binom{\ell}{k}}$$
.)*

^{*} See Appendix

Substituting the above expression into (41) and equating the coefficient of the $t^{\rm th}$ power of w, we obtain (15).

4. Remarks, Special Cases, and Examples.

A recursion formula equivalent to the identity (18) can be derived from equation (i) if K(x,y) satisfies some conditions such that integrations by parts are possible in equation (i). We have the following theorem:

Theorem 3: If $\left(\frac{\partial}{\partial y}\right)^n K(x,y) = O(e^{c \left|y\right|})$ for a fixed constant c and for all natural numbers n, then by integrations by parts we can obtain the following double recursion formula from equation (i):

$$(22) \quad \delta_{0,s}k_{r} - r! \sum_{n=r}^{s+r} (-1)^{n} {n \choose r} c_{s+r-n,n} - \sum_{m=0}^{r-1} \sum_{n=r-m-1}^{s+y-m-1} (r-m-1)! k_{m} {n \choose r-m-1}$$

$$\cdot c_{s+r-m-1-n,n} = 0$$

where

$$r_{+}(\alpha) = \sum_{n=0}^{\infty} \frac{k_{n}}{\left(\text{i}\alpha\right)^{n}} \text{, } K(x,y) = \sum_{p=0}^{\infty} \sum_{m+n=p} c_{m,n} x^{m} y^{n} \text{ being assumed.}$$

Proof: By integrating by parts N times, we obtain

$$\begin{split} & \int_{-\infty}^{\infty} \!\! K(x,y) \ e^{-i\alpha y} \mathrm{d}y = -e^{-i\alpha x} \ \sum_{n=0}^{N-1} \left(\frac{\partial}{\partial y}\right)^n \ K(x,y) \bigg]_{y=x} (i\alpha)^{-n-1} + O(\alpha^{-N-1}) \\ & \int_{-x}^{\infty} \!\! K(x,y) \ e^{i\alpha y} \mathrm{d}y = e^{-i\alpha x} \sum_{n=0}^{N-1} \left(\frac{\partial}{\partial y}\right)^n \ K(x,y) \bigg]_{y=-x} (-i\alpha)^{-n-1} + O(\alpha^{-N-1}) \,. \end{split}$$

Equating the r^{th} power of $(i\alpha)^{-1}$ in the equation (i), we obtain

$$(23) \qquad \qquad k_{\mathbf{r}} + (-1)^{\mathbf{r}-1} \left(\frac{\partial}{\partial y}\right)^{\mathbf{r}} K(\mathbf{x}, \mathbf{y}) \\ \end{bmatrix}_{\mathbf{y} = -\mathbf{x}} - \sum_{m=0}^{\mathbf{r}-1} k_{m} \left(\frac{\partial}{\partial y}\right)^{m-1} K(\mathbf{x}, \mathbf{y}) \\ \end{bmatrix}_{\mathbf{y} = \mathbf{x}} = 0 \quad (\mathbf{r} \leq \mathbf{N})$$

Putting $K(x,y) = \sum_{D=0}^{\infty} \sum_{m+n=D} c_{m,n} x^m y^n$ and equating the sth power of x

in the identity (23), we obtain the identity (22). If we assume

$$\begin{split} R(w) &= \sum_{n=0}^{\infty} \gamma_n w^n \text{ then } r_+(\alpha) = \sum_{n=0}^{\infty} \frac{\frac{k}{n}}{(i\alpha)^{n+1}} = \int_0^{\infty} R(w) e^{i\alpha w} dw \\ &= \sum_{n=0}^{\infty} \frac{(-1)^{n+1} n! \gamma_n}{(i\alpha)^{n+1}} \end{split}$$

i.e.

(24)
$$k_n = (-1)^{n+1} n! \gamma_n$$
.

Substituting the identity (24) for $k_{\rm n}$ into the identity (22) and dividing by r! in the latter, we can easily obtain the identity (18) if we make a suitable change of variables.

In Theorem 2 (or Theorem 3), computation of $C_{m,n}$ is sometimes tedious; e.g. $C_{m,n}$ (m+n \leq 5) in terms of γ_I are as follows:

[We could obtain $C_{m,n}$ in terms of k, if we used the identity (24).]

$$\begin{split} & c_{0,0} = -\gamma_{0} \\ & c_{0,1} = -\gamma_{1} + \gamma_{0}^{2} \\ & c_{0,2} = -\gamma_{2} + \frac{1}{2}\gamma_{0}^{3} \\ & c_{0,3} = -\gamma_{3} + \frac{2}{3}\gamma_{2}\gamma_{0} - \frac{1}{6}\gamma_{1}^{2} + \frac{1}{6}\chi\gamma_{0}^{2} - \frac{1}{6}\gamma_{0}^{4} \\ & c_{0,4} = -\gamma_{4} + \frac{1}{4}\gamma_{2}\gamma_{0}^{2} - \frac{1}{24}\gamma_{1}^{2}\gamma_{0} - \frac{1}{24}\gamma_{0}^{5} \\ & c_{0,5} = -\gamma_{5} + \frac{2}{5}\gamma_{4}\gamma_{0} - \frac{1}{10}\chi\gamma_{1} + \frac{1}{20}\gamma_{3}\gamma_{0}^{2} + \frac{1}{30}\gamma_{2}\gamma_{1}\gamma_{0} - \frac{1}{15}\chi_{2}\gamma_{0}^{3} \\ & - \frac{1}{120}\gamma_{1}^{3} + \frac{1}{60}\gamma_{1}^{2}\gamma_{0}^{2} - \frac{1}{120}\gamma_{1}\gamma_{0}^{4} + \frac{1}{120}\gamma_{0}^{6} \end{split}$$

$$\begin{split} & c_{1,0} = c_{0,1} = -\gamma_{1} + \gamma_{0}^{2} \\ & c_{1,1} = -2\gamma_{2} + 2\gamma_{1}\gamma_{0} - \gamma_{0}^{3} \\ & c_{1,2} = -3\gamma_{3} + \frac{1}{2}\gamma_{1}^{2} + \frac{1}{2}\gamma_{1}\gamma_{0}^{2} - \frac{1}{2}\gamma_{0}^{4} \\ & c_{1,3} = -4\gamma_{4} + 2\gamma_{3}\gamma_{0} - \frac{1}{3}\gamma_{2}\gamma_{0}^{2} + \frac{1}{6}\gamma_{1}^{2}\gamma_{0} - \frac{1}{3}\gamma_{1}\gamma_{0}^{3} + \frac{1}{6}\gamma_{0}^{5} \\ & c_{1,4} = -5\gamma_{5} + \frac{1}{2}\gamma_{3}\gamma_{1} + \frac{1}{4}\gamma_{3}\gamma_{0}^{2} - \frac{1}{6}\gamma_{2}^{2} + \frac{1}{6}\gamma_{2}\gamma_{1}\gamma_{0} - \frac{1}{6}\gamma_{2}\gamma_{0}^{2} - \frac{1}{24}\gamma_{1}^{3} \\ & - \frac{1}{24}\gamma_{1}\gamma_{0}^{4} + \frac{1}{24}\gamma_{0}^{6} \\ & c_{2,0} = -\gamma_{2} + 2\gamma_{1}\gamma_{0} - \frac{3}{2}\gamma_{0}^{3} \\ & c_{2,1} = -3\gamma_{3} + 2\gamma_{2}\gamma_{0} + \frac{3}{2}\gamma_{1}^{2} - \frac{7}{2}\gamma_{1}\gamma_{0}^{2} + \frac{3}{2}\gamma_{0}^{4} + \frac{1}{3}\gamma_{2}^{2} - \gamma_{2}\gamma_{1}\gamma_{0} + \frac{2}{3}\gamma_{2}\gamma_{0}^{3} \\ & c_{2,2} = -6\gamma_{4} + 2\gamma_{2}\gamma_{1} - \frac{1}{2}\gamma_{2}\gamma_{0}^{2} - \frac{1}{4}\gamma_{1}^{2}\gamma_{0} - \gamma_{1}\gamma_{0}^{3} + \frac{3}{4}\gamma_{0}^{5} \\ & c_{2,3} = -10\gamma_{5} + \frac{1}{4}\gamma_{4}\gamma_{0} + \gamma_{3}\gamma_{1} - \frac{3}{2}\gamma_{3}\gamma_{0}^{2} + \frac{1}{3}\gamma_{2}^{2} - \gamma_{2}\gamma_{1}\gamma_{0} + \frac{2}{3}\gamma_{2}\gamma_{0}^{3} \\ & + \frac{1}{4}\gamma_{1}^{3} - \frac{1}{2}\gamma_{1}^{2}\gamma_{0}^{2} + \frac{7}{12}\gamma_{1}\gamma_{0}^{4} - \frac{1}{4}\gamma_{0}^{6} \\ & c_{3,1} = -4\gamma_{4} + 2\gamma_{3}\gamma_{0} + 4\gamma_{2}\gamma_{1} - \frac{13}{3}\gamma_{2}\gamma_{0}^{2} - \frac{23}{6}\gamma_{1}\gamma_{0}^{2} + \frac{17}{3}\gamma_{1}\gamma_{0}^{3} - \frac{117}{6}\gamma_{0}^{5} \\ & c_{3,2} = -10\gamma_{5} + 3\gamma_{3}\gamma_{1} - \frac{3}{2}\gamma_{3}\gamma_{0}^{2} + \frac{7}{3}\gamma_{2}\gamma_{0}^{2} - \frac{23}{6}\gamma_{1}^{2}\gamma_{0} + \frac{17}{3}\gamma_{1}\gamma_{0}^{3} - \frac{117}{6}\gamma_{0}^{5} \\ & c_{3,1} = -4\gamma_{4} + 2\gamma_{3}\gamma_{0} + 4\gamma_{2}\gamma_{1} - \frac{13}{3}\gamma_{2}\gamma_{0}^{2} - \frac{23}{6}\gamma_{1}\gamma_{0}^{2} + \frac{17}{3}\gamma_{1}\gamma_{0}^{3} - \frac{117}{6}\gamma_{0}^{5} \\ & c_{3,2} = -10\gamma_{5} + 3\gamma_{3}\gamma_{1} - \frac{3}{2}\gamma_{3}\gamma_{0}^{2} + \frac{7}{3}\gamma_{2}\gamma_{0}^{2} - \frac{25}{6}\gamma_{1}\gamma_{0}^{2} + \frac{17}{3}\gamma_{1}\gamma_{0}^{3} - \frac{117}{3}\gamma_{1}\gamma_{0}^{3} - \frac{117}{6}\gamma_{0}^{5} \\ & c_{4,0} = -\gamma_{4} + 4\gamma_{3}\gamma_{0} + 2\gamma_{2}\gamma_{1} - \frac{127}{32}\gamma_{2}\gamma_{0}^{2} - \frac{27}{3}\gamma_{2}\gamma_{0}^{2} + \frac{17}{3}\gamma_{2}\gamma_{0}^{2} - \frac{1}{3}\gamma_{2}\gamma_{0}^{3} - \frac{127}{3}\gamma_{1}\gamma_{0}^{3} \\ & + \frac{13}{23}\gamma_{2}\gamma_{0}^{3} - \frac{121}{12}\gamma_{3}\gamma_{1} - \frac{129}{120}\gamma_{3}\gamma_{0}^{2} + \frac{31}{30}\gamma_{2}^{2} - \frac{239}{5}\gamma_{2}\gamma_{1}\gamma_{0} \\ & + \frac{19}{3}\gamma_{2}\gamma_{0}^{3} - \frac{12$$

 $C_{q,f}$ can also be written in determinant form as follows:

where

$$\mathbf{q}^{\mathbf{f}}_{\mathbf{1}} = \delta_{0,q} \gamma_{\mathbf{1}} + (-1)^{q} \sum_{m=0}^{q-1} (-1)^{m} \binom{t+q-m}{q-m} \mathbf{c}_{m,t+q-m} + \sum_{k=0}^{t-1} \frac{(-1)^{k} \gamma_{t-k-1}}{t \binom{t-1}{k}}$$

$$\cdot \sum_{m=0}^{q-1} \binom{k+q-m}{q-m} \mathbf{c}_{m,k+q-m} \quad (q,t=0,1,2,3....).$$

The formula (25) can be derived from the following formula (26) which is a different expression of (18), by applying Cramer's Rule for linear equations and by regarding $C_{q,k}(0 \le k \le 1)$ as unknowns.

$$(26) \quad C_{q,t} + \sum_{k=0}^{t-1} (-1)^k \frac{\gamma_{t-k-1}}{t\binom{t-1}{k}} C_{q,k} + qf_t = 0 \quad (0 \le t \le 1).$$

It is known that we can obtain $C_{0,m}$ easily by calculating (2); nevertheless it is convenient to calculate $C_{0,0}$ for $p \geq 1$ by using (10).

The integral equations ($\bar{\jmath}$) and (4) are of Volterra type if we regard K(x,y) as a given kernel and R(x) as unknown; therefore, the solution of ($\bar{\jmath}$) and (4) [as well as (1)] can be said to be an Inverse Problem of the Volterra equation. For a given K(x,w-x) in ($\bar{\jmath}$) [and in (4)], solving for R(w), we obtain a solution R(w) which in general not only depends on w but also depends on x. Consequently, the condition that R(w) is independent of x is rather strong. It implies that functions K(x,y) have a particular form. We shall discuss this situation in the following special cases:

(a) When K(x,y) = A(x)B(y), we have from the identities (23) and (24), the following:

(a,1)
$$\gamma_0 = -k_0 = -K(x,-x) = -A(x)B(-x)$$

(a,2)
$$\gamma_1 = k_1 = -\gamma_0 A(x)B(x) - A(x)B'(-x)$$

$$(a,3) \hspace{1cm} \gamma_2 = \frac{-1}{2} \; k_2 = - \; \frac{1}{2} \gamma_1 A(x) B(x) \; + \; \frac{1}{2} \gamma_0 A(x) B^{\dagger}(x) \; - \; \frac{1}{2} A(x) B^{\dagger}(-x).$$

From (a,1) and (a,2) we obtain, by eliminating A(x),

(a,4)
$$\gamma_1 B(-x) = \gamma_0^2 B(x) + \gamma_0 B'(-x)$$
.

In changing the sign of x, we have

(a,5)
$$\gamma_1 B(x) = \gamma_0^2 B(-x) + \gamma_0 B'(x)$$
.

By adding (a,4) and (a,5) and by subtracting (a,4) from (a,5) we obtain, respectively:

$$(a, 6) \qquad \gamma_0 \left\{ B^{\dagger}(x) + B^{\dagger}(-x) \right\} = -(\gamma_0^2 - \gamma_1) \left\{ B(x) + B(-x) \right\}$$

$$(a, 7) \qquad \gamma_0 \left\{ B^{\dagger}(x) - B^{\dagger}(-x) \right\} = (\gamma_0^2 + \gamma_1) \left\{ B(x) - B(-x) \right\} .$$

By putting $G(x) = \frac{1}{2} \left\{ B(x) + B(-x) \right\}, F(x) = \frac{1}{2} \left\{ B(x) - B(-x) \right\}$

and differentiating (a,6) and (a,7), we obtain

(a,8)
$$\gamma_0^2 F''(x) + (\gamma_0^4 - \gamma_1^2) F(x) = 0$$

(a,9)
$$\gamma_0^2 G''(x) + (\gamma_0^4 - \gamma_1^2)G(x) = 0$$
.

By noticing that G(x) is even and F(x) is odd, and using (a,6) and

$$(a,7)$$
, we obtain $(C = constant)$

$$G(x) = \begin{cases} \frac{a}{\gamma_{o}}x & \frac{-a}{\gamma_{o}}x \\ C(e^{\frac{a}{\gamma_{o}}} + e^{\frac{-a}{\gamma_{o}}}) & \text{if } |\gamma_{1}| > \gamma_{o}^{2} \text{ and } a = \sqrt{\gamma_{1}^{2} - \gamma_{o}^{4}} \\ \frac{a}{\gamma_{o}} \text{ix} & \frac{-a}{\gamma_{o}} \text{ix} \\ C(e^{\frac{a}{\gamma_{o}}} + e^{\frac{-a}{\gamma_{o}}}) & \text{if } |\gamma_{1}| < \gamma_{o}^{2} \text{ and } a = \sqrt{\gamma_{o}^{4} - \gamma_{1}^{2}} \end{cases}$$

$$\mathbb{F}(\mathbf{x}) = \left\{ \begin{array}{l} \frac{\gamma_1 - \gamma_0^2}{a} \frac{\frac{a}{\sigma} \mathbf{x}}{c(e^{\frac{a}{\gamma_0}} \mathbf{x} - e^{\frac{-a}{\gamma_0}} \mathbf{x})} \text{ if } |\gamma_1| > \gamma_0^2 \text{ and } a = \sqrt{\gamma_1^2 - \gamma_0^4} \\ -\frac{\gamma_0^2 - \gamma_1}{ai} \frac{\frac{a}{\gamma_0} ix}{c(e^{\frac{a}{\gamma_0}} - e^{\frac{-a}{\gamma_0}} ix} \right. \text{ if } |\gamma_1| < \gamma_0^2 \text{ and } a = \sqrt{\gamma_0^4 - \gamma_1^2} \text{.} \end{array}$$

Consequently, we obtain
$$\begin{aligned} & -\frac{a}{\gamma_o} y & \frac{a}{\gamma_o} y \\ & (26) \\ & K(x,y) = \frac{-\gamma_o B(y)}{E(-x)} \end{aligned} = \begin{cases} -\gamma_o \frac{(a+\gamma_1-\gamma_o^2)e^{-a+(a-\gamma_1+\gamma_o^2)e^{-a}}}{-\frac{a}{\gamma_o} x} & \text{if } \begin{cases} |\gamma_1| > \gamma_o^2 \\ a = \sqrt{\gamma_1^2 - \gamma_o^4} \end{cases} \\ & (a-\gamma_1+\gamma_o^2)e^{-a+(a+\gamma_1-\gamma_o^2)e^{-a}} & \text{if } \begin{cases} a = \sqrt{\gamma_1^2 - \gamma_o^4} \\ a = \sqrt{\gamma_1^2 - \gamma_o^4} \end{cases} \\ & (\alpha-\gamma_1+\gamma_o^2)e^{-a+(a-\gamma_1+\gamma_o^2)e^{-a}} & \text{if } \begin{cases} \gamma_o^2 > |\gamma_1| \\ a = \sqrt{\gamma_0^2 - \gamma_o^4} \end{cases} \\ & (\alpha-\gamma_1+\gamma_o^2)e^{-a+(a-\gamma_1+\gamma_o^2)e^{-a+(a-\gamma_1+\gamma_o^2)e^{-a}}} & \text{if } \begin{cases} \gamma_o^2 > |\gamma_1| \\ a = \sqrt{\gamma_0^4 - \gamma_o^2} \end{cases} \\ & (\alpha-\gamma_1+\gamma_o^2)e^{-a+(a-\gamma_1+\gamma_o^2)e$$

since for $\gamma_1 = \gamma_0^2$, (a, 9), (a, 7) become, respectively,

$$F^{t}(x) = 0$$
 and $G^{t}(x) = \gamma_{0}F(x)$.

It follows that F(x)=00 because F(x) is odd and hence G(x)=00 - similarly for the case $\gamma_1=-\gamma_0^{-2}$. If $\gamma_0=0$, there is only the trivial solution $K(x,y)\equiv 0$. By the relation (a,2), if $\gamma_1=0$, we have $-\gamma_0 B(x)+B^*(-x)=0$ from which we can show that $\gamma_n=0$ for n>1. From (a,3) we find that $\gamma_2=\gamma_1^{-2}/\gamma_0$ and, in general, can show that $\gamma_n=\gamma_0(\frac{\gamma_1}{\gamma_0})^n, \text{ i.e., } R(x)=\gamma_0 e^{-\frac{\gamma_1}{\gamma_0}x}$ Conversely, if $R(x)=\gamma_0 e^{-\frac{\gamma_1}{\gamma_0}x}$ then K(x,y) has exactly the form of (26); we shall show this in Example I.

(b) When K(x,y) = A(x)B(y) + C(x), we have from the identities (23) and (24), the following:

Eliminating A(x), C(x) from (b,l-3), we obtain

$$\begin{vmatrix} B(-x) & , & 1 & , & \gamma_{o} \\ B'(-x) + \gamma_{o}B(x) & , & \gamma_{o}, & \gamma_{1} \\ B''(-x) - \gamma_{o}B'(x) + \gamma_{1}B(x), & \gamma_{1}, & 2\gamma_{2} \end{vmatrix} = 0$$

which can be written as

$$(\text{b,4}) \quad \text{B"(-x)} \ - \ \gamma_{\text{o}} \text{B'(x)} \ - \frac{2 \gamma_{\text{o}} - \gamma_{\text{o}} \gamma_{\text{d}}}{\gamma_{\text{l}} - \gamma_{\text{o}}} \text{B'(-x)} - \frac{2 \gamma_{\text{o}} \gamma_{\text{c}} - \gamma_{\text{l}}^2}{\gamma_{\text{l}} - \gamma_{\text{o}}^2} \left\{ \text{B(x)} \ - \ \text{B(-x)} \right\} = 0$$

By changing the sign of x and putting K = $\frac{2\gamma_2 - \gamma_0 \gamma_1}{\gamma_1 - \gamma_0^2}$, we have

$$B''(x) - \gamma_0 B'(-x) - KB'(x) + (\gamma_1 - \gamma_0 K) \left\{ B(-x) - B(x) \right\} = 0$$

By a procedure similar to the one used in (a), we obtain

(b,5)
$$G''(x) - (\gamma + K) F'(x) = 0$$

$$\begin{array}{lll} (\text{b,6}) & F''(x) \,+\, (\,\gamma_{_{\!\!O}} \,-\, K) \,\, G^{\, t}(x) \,+\, 2(\,\gamma_{_{\!\!O}} K \,-\, \gamma_{_{\!\!O}}) \,\, F(x) \,=\, 0 \\ \\ \text{where } G(x) \,=\, \frac{1}{2} \,\left\{\,\, B(x) \,+\, B(-x)\,\,\right\} \,\, , \,\, F(x) \,=\, \frac{1}{2} \,\left\{\,\, B(x) \,-\, B(-x)\,\,\right\} \,\, . \end{array}$$

Differentiating (b,5-6) and eliminating G''(x), we have

$$H''(x) - \lambda^{2}H(x) = 0$$

where

$$H(x) = F'(x), -\lambda^2 = \gamma_0^2 - 2\gamma_1 + 2K\gamma_0 - K^2$$

It follows that if $\lambda^2 \neq 0$, we have a solution

$$B(x) = C_1 e^{\lambda x} + C_2 e^{-\lambda x} (C_1, C_2 = constants).$$

And A(x), C(x) can be computed from B(x) by solving simultaneous linear equations (b, 1-2) with respect to A(x) and C(x), i.e.,

$$A(x) = \frac{\begin{vmatrix} \gamma_{0} & 1 \\ \gamma_{1} & \gamma_{0} \end{vmatrix}}{\begin{vmatrix} -B(-x), 1 \\ -B^{\dagger}(x) + \gamma_{0}B(x), \gamma_{0} \end{vmatrix}} = \frac{\gamma_{0}^{2} - \gamma_{1}}{B^{\dagger}(-x) - \gamma_{0} \left\{ B(x) + B(-x) \right\}}$$

$$C(x) = \frac{\begin{vmatrix} -B(-x), & \gamma_{0} \\ -B(-x), & \gamma_{0} \end{vmatrix}}{\begin{vmatrix} -B(-x), & \gamma_{0} \\ -B(-x), & \gamma_{0} \end{vmatrix}} = \frac{-\gamma_{0}B^{\dagger}(-x) + \gamma_{0}^{2}B(x) + \gamma_{1}B(-x)}{B^{\dagger}(-x) - \gamma_{0} \left\{ B(x) + B(-x) \right\}}.$$

For the case γ_0 = γ_2 = 0, we have B(x) = C_1 sinh $\sqrt{2\gamma_1}$ x + C_2, hence

$$K(x,y) = A(x)B(y) + C(x) = -\sqrt{\frac{\gamma_1}{2}} \frac{\sinh\sqrt{2\gamma_1} y + \sinh\sqrt{2\gamma_1} x}{\cosh\sqrt{2\gamma_1} x} \ .$$

We also can show, in this case, $\gamma_3 = \frac{\gamma_1^2}{3!}$, etc., ...; we shall discuss the converse in Example II.

In the following we shall show some examples which can be solved by Theorems 1-3.

Example I. If
$$R(t) = be^{-ct}$$
 (b, c: constants)

$$\begin{split} r_+(\alpha) &= \int_0^\infty R(t) e^{i\alpha t} dt = \frac{-b}{i\alpha - c} \ , \qquad \text{Im}(\alpha) >\!\!> 0 \\ r_-(\alpha) &= \int_{-\infty}^\infty R(t) e^{i\alpha t} dt = \frac{b}{i\alpha - c} \ , \qquad \text{Im}(\alpha) <\!\!< 0 \ . \end{split}$$

Hence

then

$$\begin{split} \widehat{B}_{+}(\alpha) &= \frac{r_{+}(\alpha)\left[r_{-}(-\alpha)-1\right]}{l+r_{+}(\alpha)r_{-}(-\alpha)} = \frac{\left(i\alpha+c+b\right)b}{\alpha^{2}+c^{2}-b^{2}} \;, \qquad \text{Im}(\alpha) >> 0 \\ \widehat{B}_{-}(\alpha) &= \frac{-r_{-}(\alpha)\left[r_{+}(-\alpha)+1\right]}{l+r_{+}(-\alpha)r_{-}(\alpha)} = \frac{\left(i\alpha+c+b\right)b}{\alpha^{2}+c^{2}-b^{2}} \;, \qquad \text{Im}(\alpha) << 0. \end{split}$$

It follows that

$$\begin{split} \mathbb{B}(\mathbf{y}) &= \frac{1}{2\pi} \int_{-\mathbf{i}s-\infty}^{-\mathbf{i}s+\infty} \hat{\mathbb{B}}_{+}(\alpha) \mathrm{e}^{-\mathbf{i}\alpha\mathbf{y}} \mathrm{d}\alpha + \frac{1}{2\pi} \int_{-\mathbf{i}t-\infty}^{-\mathbf{i}t+\infty} \hat{\mathbb{B}}_{-}(\alpha) \mathrm{e}^{-\mathbf{i}\alpha\mathbf{y}} \mathrm{d}\alpha \\ & (s \gg 0, \, t \ll 0) \\ &= \frac{1}{2\pi} \int_{\Gamma} \frac{(\mathbf{i}\alpha + \mathrm{c} + \mathrm{b}) \mathrm{b}}{\alpha^2 + \mathrm{c}^2 - \mathrm{b}^2} \, \mathrm{e}^{-\mathbf{i}\alpha\mathbf{y}} \mathrm{d}\alpha \quad \begin{bmatrix} \mathrm{d}^2 = \mathrm{c}^2 - \mathrm{b}^2 > 0 \\ \mathrm{similarly \, for \, other} \end{bmatrix} \\ &= \frac{-1}{4\pi \mathrm{i}} \int_{\Gamma} \left\{ \frac{\mathrm{b}}{\mathrm{d}} \, \frac{\mathrm{d} - \mathrm{c} - \mathrm{b}}{(\alpha - \mathrm{i} \mathrm{d})} + \frac{\mathrm{b}}{\mathrm{d}} \, \frac{\mathrm{d} + \mathrm{c} + \mathrm{b}}{(\alpha + \mathrm{i} \mathrm{d})} \right\} \, \mathrm{e}^{-\mathrm{i}\alpha\mathbf{y}} \mathrm{d}\alpha \\ &= \frac{-\mathrm{b}}{2\mathrm{d}} \left\{ (\mathrm{d} - \mathrm{c} - \mathrm{b}) \, \mathrm{e}^{\mathrm{d}\mathbf{y}} + (\mathrm{d} + \mathrm{c} + \mathrm{b}) \, \mathrm{e}^{-\mathrm{d}\mathbf{y}} \right\} \end{split}$$

Trying the form K(x,y) = A(x)B(y) for a solution of equation (1), we obtain exactly formula (26) if we put $\gamma_0 = b$, $\gamma_1 = -bc$.

shown above.

of two parallel lines as

Substituting
$$K(x,y) = \frac{-b}{2d} A(x) \left\{ (d-c-b)e^{dy} + (d+c+b)e^{-dy} \right\}$$
 into

(1), and solving for A(x), we obtain (26). In the above case, of course, we could assume $K(x,y)=\sum_{n=0}^{\infty}A_{n}(y)x^{n}$ and apply Theorem 1 for

computation of $A_n(y)$, step by step; nevertheless, the computation is not very simple. For example, we have

$$\begin{split} A_{o}(y) &= B(y) = \frac{-b}{2d} \left\{ (d - c - b)e^{dy} + (d + c + b)e^{-dy} \right\} \\ R^{1}(x,t) &= \frac{1}{x} \left\{ b \int_{-t}^{x} e^{-c(t + y)} A_{o}(y) dy + A_{o}(t) + be^{-c(t + x)} \right\} \\ &= \frac{b}{x} \left\{ \int_{0}^{x} e^{-c(t + y)} A_{o}(y) dy + e^{-c(t + x)} - e^{-ct} \right\} \end{split}$$

since
$$\int_{-t}^{0} be^{-c(t+y)} A_{o}(y) dy + A_{o}(t) + be^{-ct} = 0$$
.

It follows that

$$R^{1}(0,t) = b \left\{ e^{-c(t+x)} A_{0}(x) \Big|_{x=0} + (-c)e^{-ct} \right\}$$

= $-b(b+c) e^{-ct}$.

Therefore

$$\begin{split} &r_{+}^{-1}(\alpha) = \int_{0}^{\infty} R^{1}(0,t) e^{i\alpha t} dt = \frac{b(b+c)}{i\alpha-c} \;, \qquad \text{Im}(\alpha) >> 0 \\ &r_{-}^{-1}(\alpha) = \int_{-\infty}^{0} R^{1}(0,t) e^{i\alpha t} dt = \frac{-b(b+c)}{i\alpha-c} \;, \qquad \text{Im}(\alpha) << 0 \;. \end{split}$$

Consequently

$$\begin{split} A_{\underline{l}}(y) &= \frac{-\dot{o}(b+c)}{2\pi} \int_{\Gamma} \frac{1\alpha + b + c}{\alpha^2 + c^2 - b^2} \ e^{-iy\alpha} d\alpha \ \left[\begin{array}{c} \Gamma \ \text{is the same contour} \\ \text{as in p. 20.} \end{array} \right] \\ &= \frac{b(b+c)}{2d} \left\{ \ (d-c-b)e^{dy} + \ (d+c+b) \ e^{-dy} \right\} \,. \end{split}$$

Similarly we can compute $A_2(y)$ as follows:

$$\begin{split} A_{\underline{1}}(y) &= \frac{b(b+c)}{2d} \, \left\{ \, (d-c-b)e^{dy} \, + \, (d+c+b)e^{-dy} \, \right\} = \, -(b+c)A_{\underline{0}}(y) \\ R^2(x,t) &= \frac{1}{x} \, \left\{ \, b \, \int_{-t}^{\, x} e^{-c(t+y)} A_{\underline{1}}(y) \, dy \, + \, A_{\underline{1}}(t) \, + \, R^{\underline{1}}(x,t) \, \right\} \\ &= \frac{1}{x} \, \left\{ \, b \, \int_{0}^{\, x} e^{-c(t+y)} A_{\underline{1}}(y) \, dy \, + \, R^{\underline{1}}(x,t) \, - \, R^{\underline{1}}(0,t) \, \right\} \, . \end{split}$$

It follows that

$$\begin{split} R^2(0,t) &= \left. b e^{-ct} A_1(0) + \frac{\partial}{\partial x} R^1(x,t) \right|_{x=0} \\ &= \left. b^2(b+c) e^{-ct} + \frac{1}{2} b(b+c)^2 e^{-ct} = \frac{1}{2} b(b+c) (3b+c) e^{-ct} \right. \end{split}$$

By means of the same procedure as we used in computing $\boldsymbol{A}_{\!\!\!\!\!\!\!\boldsymbol{\gamma}}\left(\boldsymbol{y}\right),$ we obtain

$$\begin{split} \mathbf{A}_{2}(\mathbf{y}) &= \frac{1}{4d} \, b(\mathbf{b} + \mathbf{c}) \, (\vec{\jmath} \mathbf{b} + \mathbf{c}) \, \left\{ (\mathbf{d} - \mathbf{c} - \mathbf{b}) \, \mathbf{e}^{\vec{\mathbf{d}} \vec{\mathbf{y}}} + \, (\mathbf{d} + \mathbf{c} + \mathbf{b}) \, \mathbf{e}^{-\vec{\mathbf{d}} \vec{\mathbf{y}}} \right\} \\ &= \frac{1}{2} (\mathbf{b} + \mathbf{c}) \, (\vec{\jmath} \mathbf{b} + \mathbf{c}) \, \mathbf{A}_{0}(\mathbf{y}) \, . \end{split}$$

Similarly, we obtain

$$\begin{split} &A_{3}(y) = -\frac{1}{3!} \; (b+c)^{2} (llb+c) A_{o}(y) \\ &A_{l_{4}}(y) = \frac{1}{l_{1}!} \; (b+c)^{2} (57b^{2} + 38bc + c^{2}) A_{o}(y) \\ &A_{5}(y) = -\frac{1}{5!} \; (b+c)^{3} (36lb^{2} + ll8bc + c^{2}) \; A_{o}(y) \quad \text{and so on.} \end{split}$$

Hence

$$K(\mathbf{x}, y) = \sum_{n=0}^{\infty} A_n(y) x^n = A_0(y) \left\{ 1 - (b+c)x + \frac{1}{2}(b+c)(3b+c)x^2 - \dots \right\}$$

from which use of the form K(x,y) = A(x)B(y) is justified.

K(x,y) can be obtained from Theorem ? (or Theorem) or follows:

Since R(t) = be^-et = b
$$\sum_{n=0}^{\infty} \frac{(-1)^n e^n}{n!}$$
 tⁿ , we have $\gamma_n = \frac{(-1)^n e^n}{n!}$ b .

Using the general formula for $C_{m,n}$ which we computed before, we obtain:

$$\begin{split} & \text{C}_{0,0} = -b \\ & \text{C}_{0,1} = b(\text{c}+b) \\ & \text{C}_{0,2} = \frac{b}{2}(-\text{c}^2 + \text{b}^2) = -\frac{1}{2}\text{bd}^2 \\ & \text{C}_{0,3} = \frac{b}{2}(\text{c}^3 + \text{bc}^2 - \text{b}^2 \text{c} - \text{b}^3) = \frac{\text{bd}^2}{3!} (\text{c}+\text{b}) \\ & \text{C}_{0,3} = \frac{b}{2!}(\text{c}^4 - 2\text{b}^2 \text{c}^2 + \text{b}^4) = \frac{-\text{bd}^4}{4!} \\ & \text{C}_{0,4} = \frac{-b}{2!}(\text{c}^4 - 2\text{b}^2 \text{c}^2 + \text{b}^4) = \frac{-\text{bd}^4}{4!} \\ & \text{C}_{0,5} = \frac{b}{120}(\text{c}^5 + \text{bc}^4 - 2\text{bc}^3 - 2\text{b}^2 \text{c}^2 + \text{b}^3 \text{c} + \text{b}^4) \\ & = \frac{b\text{d}^2}{5} (\text{c}+\text{b})^3 \\ & \text{C}_{2,0} = \frac{-b}{2}(\text{c}^2 + 4\text{bc} + 3\text{b}^2) = \frac{-b}{2}(\text{c}+\text{b})(\text{c} + 3\text{b}) \\ & = \frac{1}{2}(\text{c}+\text{b})(\text{c} + 3\text{b})\text{C}_{0,0} \\ & \text{C}_{2,1} = \frac{b}{2}(\text{c}^3 + 5\text{bc}^2 - 7\text{b}^2 \text{c} + 3\text{b}^3) \\ & = \frac{b}{2}(\text{c}+\text{b})^2(\text{c} + 3\text{b}) = \frac{1}{2}(\text{c}+\text{b})(\text{c} + 3\text{b})\text{C}_{0,1} \\ & \text{C}_{2,2} = \frac{-b}{4}(\text{c}^4 + 4\text{bc}^3 + 2\text{b}^2 \text{c}^2 - 4\text{b}^3 \text{c} - 3\text{b}^2) \\ & = -\frac{b}{4}\text{d}^2(\text{c}+\text{b})(\text{c} + 3\text{b}) = \frac{1}{2}(\text{c}+\text{b})(\text{c} + 3\text{b})\text{C}_{0,2} \\ & \text{C}_{2,3} = \frac{b}{12}(\text{c}^5 + 5\text{bc}^4 + 6\text{b}^2 \text{c}^3 - 2\text{b}^3 \text{c}^2 - 7\text{b}^4 \text{c} - 3\text{b}^5) \\ & = \frac{b}{12}\text{d}^2(\text{c}+\text{b})^2(\text{c}^4 + 3\text{bb}\text{c}^3 + 134\text{b}^2 \text{c}^2 + 152\text{b}^3 \text{c} + 57\text{b}^4) \\ & = -\frac{b}{44}(\text{c}+\text{b})^2(\text{c}^2 + 38\text{bc} + 57\text{b}^2) \\ & = \frac{1}{44}(\text{c}+\text{b})^2(\text{c}^2 + 38\text{bc} + 57\text{b}^2) \\ & = \frac{1}{44}(\text{c}$$

$$\begin{split} c_{i_{1},1} &= \frac{b}{2h}(c^{5} + 41bc^{4} + 174b^{2}c^{3} + 255b^{2}c^{2} \\ &\quad + 209b^{4}c + 57b^{5}) \\ &= \frac{b}{4!}(c+b)^{3}(c^{2} + 38bc + 57b^{2}) \\ &= \frac{1}{h!}(c+b)^{2}(c^{2} + 38bc + 57b^{2})c_{1,0} \end{split}$$

$$\begin{split} & c_{1,0} = b(c+b) = -(c+b)c_{0,0} \\ & c_{1,1} = -b(c^2 + 2bc + b^2) + -b(c+b)^2 \\ & = -(c+b)c_{0,1} \\ & c_{1,2} = \frac{b}{2}(c^3 + bc^2 - b^2c - b^3) = \frac{bd^2}{2}(c+b) \\ & = -(c+b)c_{0,2} \\ & c_{1,3} = \frac{-b}{6}(c^4 + 2c^3b - 2cb^3 - b^4) \\ & = \frac{-bd^2}{3!}(c+b)^2 = -(c+b)c_{0,3} \\ & c_{1,4} = \frac{b}{24}(c^5 + bc^4 - 2b^2c^3 - 2b^3c^2 + b^4c + b^5) \\ & = \frac{bd^4}{4!}(c+b) = -(c+b)c_{0,4} \\ & c_{3,0} = \frac{b}{6}(c^3 + 13bc^2 + 23b^2c + 11b^2) \\ & = \frac{b}{3!}(c+b)^2(c+11b) = \frac{-1}{3!}(c+b)^2(c+11b)c_{0,0} \\ & c_{3,1} = \frac{-b}{6}(c^4 + 14bc + 36b^2c^2 + 34b^3c + 11b^4) \\ & = \frac{-b}{3!}(c+b)^3(c+11b) = \frac{-1}{3!}(c+b)^2(c+11b)c_{0,1} \\ & c_{3,2} = \frac{b}{12}(c^5 + 13bc^4 + 22b^2c^3 - 2b^3c^2 - 23b^4c - 11b^5) \\ & = \frac{bd^2}{12}(c+1)^2(c+11b) = \frac{-1}{3!}(c+b)^2(c+11b)c_{0,2} \\ & c_{5,0} = \frac{b}{120}(c^5 + 121bc^4 + 718b^2c^3 + 1438b^3c^2 \\ & + 1201b^4c + 361b^5) \\ & = \frac{b}{5!}(c+b)^3(c^2 + 118bc + 361b^2)c_{0,0} \\ & = \frac{-1}{5!}(c+b)^3(c^2 + 118bc + 361b^2)c_{0,0} \\ \end{split}$$

These coefficients $C_{m,n}$ are exactly the coefficients of powers $x^m y^n$ in

the expansion of
$$K(x,y) = -b \frac{(d-c-b)e^{dy} + (d+c+b)e^{-dy}}{(d+c+b)e^{dy} + (d-c-b)e^{-dy}}$$
 (if $d^2 = c^2 - b^2 > 0$, similarly for $d^2 \le 0$.)

Example II. If
$$R(t) = \sqrt{b} \sinh \sqrt{b} t = \frac{\sqrt{b}}{2} (e^{\sqrt{b}t} - e^{-\sqrt{b}t})$$
 then
$$K(x,y) = -\sqrt{\frac{b}{2}} \frac{\sinh \sqrt{2b} x + \sinh \sqrt{2b} y}{\cosh \sqrt{2b} x}.$$

To show this, applying Theorem 1, we have

$$\begin{split} \mathbf{r}_{+}(\alpha) &= \int_{0}^{\infty} \mathbf{R}(\mathbf{t}) \ \mathrm{e}^{\mathrm{i}\alpha t} \mathrm{d}\mathbf{t} = \frac{-b}{\alpha^{2} + b} \ , \quad \mathrm{Im}(\alpha) \gg 0 \\ \\ \mathbf{r}_{-}(\alpha) &= \int_{-\infty}^{0} \mathbf{R}(\mathbf{t}) \ \mathrm{e}^{\mathrm{i}\alpha t} \mathrm{d}\mathbf{t} = \frac{b}{\alpha^{2} + b} \ , \quad \mathrm{Im}(\alpha) \gg 0 \ . \end{split}$$

From the formula (7), we obtain

$$\begin{split} \hat{K}_{+}(0,\alpha) &= \frac{b}{\alpha^{2} + 2b} \;, \qquad \text{Im}(\alpha) \gg 0 \\ \hat{K}_{-}(0,\alpha) &= \frac{-b}{\alpha^{2} + 2b} \;, \qquad \text{Im}(\alpha) \ll 0 \;. \end{split}$$

Hence

$$\begin{split} A_{_{\scriptsize O}}(y) &= K(0,y) = \frac{-1}{2\pi} \int_{\Gamma} \frac{be^{-iy\alpha}}{\alpha^2 + 2b} \, \mathrm{d}\alpha \, \left[\begin{array}{c} \Gamma \text{ is the same contour} \\ \text{as in Example I.} \end{array} \right] \\ &= \frac{-b}{2\,\sqrt{2}}\, \frac{1}{2\pi i} \, \int_{\Gamma} \, \left\{ \frac{1}{\alpha - i\,\sqrt{2b}} - \frac{1}{\alpha + i\,\sqrt{2b}} \right\} \, \mathrm{e}^{-iy\alpha} \mathrm{d}\alpha \\ &= -\sqrt{\frac{b}{2}} \, \mathrm{sinh}\, \sqrt{2b} \, y \, . \end{split}$$

Similarly, we obtain

$$A_1(y) = -b$$

^{*} For b = 1, see [1], p. 19-21.

$$\begin{split} & A_2(y) \, = \, b \sqrt{\frac{b}{2}} \quad \sinh \sqrt{2b} \ y \\ & A_3(y) \, = \, \frac{2}{3} \ b^2 \\ & A_4(y) \, = \, \frac{-5}{6} \, \sqrt{\frac{b}{2}} \ b^2 \ \sinh \ \sqrt{2b} \ y \\ & A_5(y) \, = \, \frac{-8}{15} \ b^3 \qquad \qquad \text{etc.} \end{split}$$

Consequently,

$$\begin{split} \mathtt{K}(\mathtt{x},\mathtt{y}) &= \sum_{n=0}^{\infty} \ \mathtt{A}_{n}(\mathtt{y}) \, \mathtt{x}^{n} \, = \, - \, \sqrt{\frac{\mathtt{b}}{2}} \quad \mathrm{sinh} \quad \sqrt{2\mathtt{b}} \ \mathtt{y}(\mathtt{1} \text{-} \mathtt{6x} + \frac{5}{6} \, \mathtt{b}^{\, 2} \mathtt{x} + \dots) \\ &+ \, \left(-\mathtt{bx} + \frac{2}{3} \mathtt{b}^{\, 2} \mathtt{x}^{\, 3} - \frac{8}{15} \mathtt{b}^{\, 3} \mathtt{x}^{\, 5} + \, \dots \right) \end{split}$$

from which it is reasonable for us to assume that K(x,y) has the form $-\sqrt{\frac{b}{2}}$ A(x)sin $\sqrt{2b}$ y + C(x); substituting the latter into (1), we obtain A(x) = sech $\sqrt{2b}$ x , C(x) = $-\sqrt{\frac{b}{2}}$ tanh $\sqrt{2b}$ x .

It also can be seen from Theorem 2 (or Theorem 3) that, in this case, we have γ_{2n} = 0, γ_{2n+1} = $\frac{b^{n+1}}{(2n+1)!}$; hence

from which we could also guess that K(x,y) is of the form

$$A(x)B(y) + C(x) \quad \text{with} \quad B(y) = -\sqrt{\frac{b}{2}} \quad \sinh \sqrt{2b} \ y, \ C(x) = -\sqrt{\frac{b}{2}} \quad \tanh \sqrt{2b} \ x \ .$$
 Substituting these into (1), we obtain $A(x) = \operatorname{sech} \sqrt{2b} \ x$.

From the above example, it is known that Theorem 2 is useful in determining the form of K(x,y) for a given R(t); similarly, for a given $r_+(\alpha)$, Theorem 3 is also useful in determining the form of K(x,y). As in Example II, it is found that if R(t) is an odd function (i.e. $\gamma_{2n}=0$), then $C_{m,n}=0$ for m+n is even *; and hence K(x,x) is also an odd function. Consequently, V(x)=2 $\frac{d}{dx}$ K(x,x) is an even function. When R(t) is even we cannot conclude very much as in the following example.

Example III If
$$R(z) = e^{-cz^2}$$
, then $\gamma_{2m} = \frac{(-1)^m c^m}{m}$, $\gamma_{2m+1} = 0$.

From Theorem 2, we then obtain $C_{p,q}(p+q \le 5)$ as follows:

$$C_{5,0} = \frac{127}{30} c^2 - \frac{223}{30} c + \frac{361}{120}$$
 and so on.

^{*}See Appendix

And

$$\begin{split} K(x,x) &= -1 + 2x + 2(2c - 1)x^2 + \frac{16}{3}(c + \frac{1}{2})x^3 + (-8c^2 + \frac{28}{3}c - \frac{10}{3})x^{\frac{1}{4}} \\ &+ (\frac{64}{5}c^2 - \frac{203}{15}c + \frac{64}{15})x^5 + \dots \end{split}$$

Example IV

If
$$R(z) = ze^{-cz^2}$$
, then $\gamma_{2m} = 0$, $\gamma_{2m+1} = \frac{(-1)^m c^m}{m!}$.

And we obtain $C_{m,n} = 0$ if m+n is even.

$$c_{3,2} = -5c^2 - 3c - \frac{5}{12}$$
, $c_{4,1} = -\frac{5}{2}c^2 - \frac{11}{2}c - \frac{41}{24}$, $c_{5,0} = -\frac{c^2}{2} - \frac{29}{10}c - \frac{121}{120}$;

hence
$$K(x,x) = -2x + (6c + \frac{16}{6})x^{3} - (16c^{2} + \frac{64}{5}c + \frac{144}{15})x^{5} + \dots$$

$$\text{ and } \mathbb{V}(\mathbb{X}) \ = \ 2 \ \frac{\mathrm{d}}{\mathrm{d} x} \ \mathbb{K}(\mathbb{X}, \mathbb{X}) \ = \ -4 \ + \ 6(6 c \ + \ \frac{16}{6}) x^2 - \ 10(16 c^2 + \ \frac{64}{5} \ c \ + \ \frac{144}{15}) x^4 \ + \ \dots$$

(i) holds since

$$\begin{split} \int_0^\infty \int_0^w R(w-u)K(x,x-u)\,\mathrm{d}u & \ \mathrm{e}^{\mathrm{i} c t w} \mathrm{d}w = \int_0^\infty K(x,x-u)\,\mathrm{e}^{\mathrm{i} t t t u} \mathrm{d}u \int_0^\infty R(w-u)\,\mathrm{e}^{\mathrm{i} c t (w-u)}_{uw} \\ & = \ \mathbf{r}_+(\alpha) \int_0^\infty K(x,x-u)\,\mathrm{e}^{\mathrm{i} t t u u} \mathrm{d}u = \ \mathbf{r}_+(\alpha) \int_{-\infty}^X K(x,y)\,\mathrm{e}^{-\mathrm{i} t t y} \mathrm{d}y + \mathrm{e}^{\mathrm{i} t t x} \\ & \int_0^\infty K(x,w-x)\,\mathrm{e}^{\mathrm{i} t t w} \mathrm{d}w = \int_X^\infty K(x,y)\,\mathrm{e}^{\mathrm{i} t t y} \mathrm{d}y + \mathrm{e}^{\mathrm{i} t t x} \end{split}$$

Similarly, (ii) holds since

$$\begin{split} \int_{-\infty}^{\circ} \int_{0}^{w} R(w-u)K(x,x-u) \, du \ e^{i\alpha w} dw &= -\int_{-\infty}^{\circ} K(x,x-u) e^{i\alpha u} du \ -\int_{-\infty}^{u} R(w-u) e^{i\alpha (w-u)} \, dw \\ &= -\mathbf{r}_{-}(\alpha) \int_{-\infty}^{\circ} K(x,x-u) e^{i\alpha u} du = -\mathbf{r}_{-}(\alpha) \int_{x}^{\infty} K(x,y) e^{-i\alpha y} dy \cdot e^{i\alpha x} \\ &\int_{-\infty}^{\circ} K(x,w-x) e^{i\alpha w} dw = \int_{-\infty}^{-x} K(x,y) e^{i\alpha y} dy \cdot e^{i\alpha x} \ . \end{split}$$

The identity (*) can be derived as follows:

$$\left(1-x \right)^b = \sum_{s=0}^b \left(-1 \right)^s \binom{b}{s} \ x^s \ , \quad \left(1-x \right)^{-a} = \sum_{r=0}^\infty \binom{a+r-1}{a-1} \ x^r$$

and

$$(1-x)^{b-a} = \sum_{t=0}^{b-a} (-1)^t \binom{b-a}{t} x^t \quad \text{(we assume } b \ge a \text{, b, a are natural numbers)}$$

It follows that by equating the t^{th} power of $(1-x)^{b-a}$ and the product of $(1-x)^b$ and $(1-x)^{-a}$, we obtain

$$(-1)^{t} \begin{pmatrix} b-a \\ t \end{pmatrix} = \sum_{s+r=t} (-1)^{s} \begin{pmatrix} b \\ s \end{pmatrix} \begin{pmatrix} a+r-1 \\ a-1 \end{pmatrix} = \sum_{s=0}^{t} (-1)^{s} \begin{pmatrix} b \\ s \end{pmatrix} \begin{pmatrix} a-1+t-s \\ a-1 \end{pmatrix} ,$$

putting a = k+1, $b = \ell+1$, $t = \ell-k$, we obtain

$$(**) \qquad (-1)^{\ell-k} = \sum_{s=0}^{\ell-k} (-1)^{s} {\ell+1 \choose s} {\ell-s \choose k} .$$

And since

by substituting the righthand side of the above into (**) and dividing by $(-1)^{l-k}$, we obtain the identity (*) easily.

To prove
$$C_{m,n} = 0$$
 for $m+n = even$, if $\gamma_{2m} = 0$:

From (18), it is easily shown that $\boldsymbol{C}_{m,n}$ has at most the following terms:

$$\gamma_{m+n}, \gamma_{m+n-1}\gamma_{o}, \gamma_{m+n-2}\gamma_{1}, \gamma_{m+n-2}\gamma_{o}^{2}, \dots$$

These are isomorphic to all partitions of m+n+1 into natural numbers such that ℓ corresponds to $\gamma_{\ell-1}$; e.g., $(m+n+1) \leftrightarrow \gamma_{m+n}; (m+n-1,1,1) \leftrightarrow \gamma_{m+n-2} \gamma_0^2$ etc. Therefore, if m + n is even, then m + n + 1 is odd; hence all their partitions are of the form (p_1, p_2, \ldots, p_i) such that one of $p_i (1 \le i \le m+n+1)$ is odd. Otherwise the sum $\sum p_i$ is not an odd number since sums of even numbers are even. It follows that $C_{m,n} = 0$ (if m+n = even) since each of their terms has at least one factor $\gamma_{p_1-1} = 0$ which corresponds to odd p_i .

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